

Stability Conditions for Cluster Synchronization in Networks of Heterogeneous Kuramoto Oscillators

Tommaso Menara , *Student Member, IEEE*, Giacomo Baggio , Danielle S. Bassett ,
and Fabio Pasqualetti , *Member, IEEE*

Abstract—In this paper, we study cluster synchronization in networks of oscillators with heterogeneous Kuramoto dynamics, where multiple groups of oscillators with identical phases coexist in a connected network. Cluster synchronization is at the basis of several biological and technological processes; yet, the underlying mechanisms to enable the cluster synchronization of Kuramoto oscillators have remained elusive. In this paper, we derive quantitative conditions on the network weights, cluster configuration, and oscillators' natural frequency that ensure the asymptotic stability of the cluster synchronization manifold; that is, the ability to recover the desired cluster synchronization configuration following a perturbation of the oscillators' states. Qualitatively, our results show that cluster synchronization is stable when the intracluster coupling is sufficiently stronger than the intercluster coupling, the natural frequencies of the oscillators in distinct clusters are sufficiently different, or, in the case of two clusters, when the intracluster dynamics is homogeneous. We validate the effectiveness of our theoretical results via numerical studies.

Index Terms—Biological networks, Kuramoto oscillators, nonlinear systems, networks of autonomous agents.

I. INTRODUCTION

SYNCHRONIZATION refers broadly to patterns of coordinated activity that arise spontaneously or by design in several natural and man-made systems [1]–[3]. Examples include coherent firing of neuronal populations in the brain [4], coordinated flashing of fireflies [5], flocking of birds [6], exchange of signals in wireless networks [7], consensus in multiagent systems [8], and power generation in the smart grid [9]. Synchronization enables complex functions: while some systems

require complete (or full) synchronization among all components in order to function properly, others rely on cluster (or partial) synchronization, where different groups exhibit different, yet synchronized, internal behaviors [10].

While studies of full synchronization are numerous and have generated rich literature (e.g., see [11]–[13]), conditions explaining the onset of cluster synchronization and its properties are less well understood. Such conditions are necessary for the analysis and, more important, the control of synchronized activity across biological [14]–[16] and technological [17] systems. For instance, a deeper understanding of the mechanisms enabling cluster synchronization might not only shed light on the nature of the healthy human brain [18], but also enable and guide targeted interventions for patients with neurological disorders, such as epilepsy [19] and Parkinson's disease [20].

We study cluster synchronization in networks of oscillators with Kuramoto dynamics [21], which, despite their apparent simplicity, are particularly suited to represent complex synchronization phenomena in neural systems [22], as well as in many other natural and technological systems [9]. Although our study and modeling choices are guided by the practical need to understand and control patterns of synchronized functional activity in the human brain, as they naturally arise in healthy and diseased populations [23], [24], in this paper, we focus on developing the mathematical foundations of a quantitative approach to the analysis and control of cluster synchronization in a weighted network of Kuramoto oscillators. In particular, we derive conditions on the oscillators' coupling and their natural frequencies that guarantee the stability of an arbitrary cluster configuration.

Related work: Cluster synchronization, where multiple synchronized groups of oscillators coexist in a connected network, is an exciting phenomenon that has attracted the attention of the physics, dynamical systems, and controls communities, among others. The existing work on this topic has shown that cluster-synchronized states can be linked to the existence of certain network symmetries [25]–[29] or symmetries in the nodes' dynamics [30]. More recently, in [31] and [32], the stability of cluster states corresponding to network symmetries is addressed with the master stability function approach [33]. In contrast to this previous work, [34] combines network symmetries with contraction analysis to study the stability of synchronized states. Further studies relating contraction properties and cluster synchronization are conducted in [35] and [36]. Finally, control

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T. Menara, G. Baggio, and F. Pasqualetti are with the Department of Mechanical Engineering, University of California at Riverside, Riverside, CA 92521 USA (e-mail: tomenara@engr.ucr.edu; gbaggio@engr.ucr.edu; fabiopas@engr.ucr.edu).

D. S. Bassett is with the Department of Bioengineering, the Department of Electrical and Systems Engineering, the Department of Physics and Astronomy, the Department of Psychiatry, and the Department of Neurology, University of Pennsylvania, Philadelphia, PA 19104 USA (e-mail: dsb@seas.upenn.edu).

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algorithms for cluster synchronization are developed in [37] and [38]. To the best of our knowledge, however, the aforementioned studies are not applicable to oscillators with Kuramoto dynamics, which we study in this paper.

A few papers have studied cluster synchronization of Kuramoto oscillators. Specifically, in [39] and [40], the authors provide invariance conditions for an approximate definition of cluster synchronization and for particular types of networks. Invariance of exact cluster synchronization, which is the notion used in this paper, is also studied in [41] and [42]. Stability of exact cluster synchronization is investigated in [43], where, however, only the restrictive case of two clusters for identical Kuramoto oscillators with inertia is considered, and in [44], where only implicit and numerical stability conditions are provided. To the best of our knowledge, our work presents the first explicit analytical conditions for the (local) stability of the cluster synchronization manifold in sparse and weighted networks of heterogeneous Kuramoto oscillators.

Paper contribution: The main contribution of this paper is to characterize conditions for the stability of cluster synchronization in networks of oscillators with Kuramoto dynamics. We consider a notion of exact cluster synchronization, where the phases of the oscillators within each cluster remain equal to each other over time, and different from the phases of the oscillators in the other clusters. We derive three conditions. First, we show that the cluster synchronization manifold is locally exponentially stable when the intracluster coupling is sufficiently stronger than the intercluster coupling. We quantify this trade-off using the perturbation theory for dynamical systems together with the invariance properties of cluster synchronization. Second, through a Lyapunov argument, we show that the cluster synchronization manifold is locally exponentially stable when the natural frequencies of the oscillators in disjoint clusters are sufficiently different (in their limit to infinity). Third, we focus on the case of two clusters, and provide a quantitative condition on the network weights and oscillators' natural frequency for the stability of the cluster synchronization manifold. This analysis shows that the asymptotic stability of the cluster synchronization manifold is guaranteed for weak intercluster weights, sufficiently different natural frequencies, or even homogeneous intracluster configurations.

As minor contributions, we provide examples showing that network symmetries are not necessary for cluster synchronization of Kuramoto oscillators, and a sufficient condition guaranteeing the absence of stable synchronization submanifolds.

Paper organization: The rest of this paper is organized as follows. Section II contains the problem setup and some preliminary notions. Section III contains the main results; that is, the conditions for the stability of the cluster synchronization manifold in Kuramoto networks. Finally, Section IV concludes this paper, and the Appendix contains the proofs of our results.

Mathematical notation: The set $\mathbb{R}_{>0}$ (resp., $\mathbb{R}_{<0}$) denotes the positive (resp., negative) real numbers, whereas the sets \mathbb{S}^1 and \mathbb{T}^n denote the unit circle and the n -dimensional torus, respectively. The vector of all ones is represented by $\mathbf{1}$. We let $O(f)$ denote the order of the function f . Further, we denote a positive (resp., negative) definite matrix A with $A \succ 0$ (resp., $A \prec 0$).

We indicate the smallest (resp. largest) eigenvalue of a symmetric matrix with $\lambda_{\min}(\cdot)$ (resp., $\lambda_{\max}(\cdot)$). A (block-)diagonal matrix is represented by $(\text{blk-})\text{diag}(\cdot)$. We let $\|\cdot\|$ denote the ℓ^2 -norm, and $i = \sqrt{-1}$. Finally, A^\dagger represents the Moore–Penrose pseudoinverse of the matrix A .

II. PROBLEM SETUP AND PRELIMINARY NOTIONS

In this paper, we characterize the stability properties of certain synchronized trajectories arising in networks of oscillators with Kuramoto dynamics. To this aim, let $\mathcal{G} = (\mathcal{V}, \mathcal{E})$ be the connected and weighted graph representing the network of oscillators, where $\mathcal{V} = \{1, \dots, n\}$ and $\mathcal{E} \subseteq \mathcal{V} \times \mathcal{V}$ represent the oscillators, or nodes, and their interconnection edges, respectively. Let $A = [a_{ij}]$ be the weighted adjacency matrix of \mathcal{G} , where $a_{ij} \in \mathbb{R}_{>0}$ is the weight of the edge $(i, j) \in \mathcal{E}$, and $a_{ij} = 0$ when $(i, j) \notin \mathcal{E}$. The dynamics of the i th oscillator is

$$\dot{\theta}_i = \omega_i + \sum_{j \neq i} a_{ij} \sin(\theta_j - \theta_i) \quad (1)$$

where $\omega_i \in \mathbb{R}_{>0}$ and $\theta_i \in \mathbb{S}^1$ denote the natural frequency and the phase of the i th oscillator. Unless specified differently, we assume that the edge weights are symmetric as follows:

(A1) The network adjacency matrix satisfies $A = A^\top$.

Assumption (A1) is typical in the study of (cluster) synchronization in networks of Kuramoto oscillators, e.g., see [45]–[47], as it facilitates the derivation of stability results. While relaxing this assumption is beyond the scope of this paper, we will discuss how our stability results can also be applied to study cluster synchronization with asymmetric network weights (see Remark 5). Finally, since the diagonal entries of the adjacency matrix A do not contribute to the dynamics in (1), we assume that \mathcal{G} does not contain self-loops.

A network exhibits cluster synchronization when the oscillators can be partitioned so that the phases of the oscillators in each cluster evolve identically. To be precise, let $\mathcal{P} = \{\mathcal{P}_1, \dots, \mathcal{P}_m\}$, with $m > 1$, be a partition of \mathcal{V} , where $\bigcup_{i=1}^m \mathcal{P}_i = \mathcal{V}$ and $\mathcal{P}_i \cap \mathcal{P}_j = \emptyset$ if $i \neq j$. Define the *cluster synchronization manifold* associated with the partition \mathcal{P} as

$$\mathcal{S}_{\mathcal{P}} = \{\theta \in \mathbb{T}^n : \theta_i = \theta_j \text{ for all } i, j \in \mathcal{P}_k, k = 1, \dots, m\}.$$

Then, the network is cluster-synchronized with partition \mathcal{P} when the phases of the oscillators belong to $\mathcal{S}_{\mathcal{P}}$ at all times.

In this paper, we characterize conditions on the network weights and the oscillators' natural frequency that guarantee *local exponential stability* of the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$, for a given partition \mathcal{P} .¹ In order to study the stability of the cluster synchronization manifold, we assume $\mathcal{S}_{\mathcal{P}}$ to be invariant [48, ch. 3].² In particular, following [42], invariance of $\mathcal{S}_{\mathcal{P}}$ is guaranteed by the following conditions:

¹Loosely speaking, the manifold $\mathcal{S}_{\mathcal{P}}$ is locally exponentially stable if θ converge to $\mathcal{S}_{\mathcal{P}}$ exponentially fast when $\theta(0)$ is sufficiently close to $\mathcal{S}_{\mathcal{P}}$.

²The manifold $\mathcal{S}_{\mathcal{P}}$ is invariant if $\theta(0) \in \mathcal{S}_{\mathcal{P}}$ implies $\theta \in \mathcal{S}_{\mathcal{P}}$ at all times.

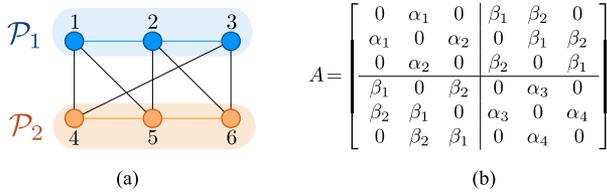


Fig. 1. (a) Network of six oscillators with adjacency matrix as in (b). In this network, the partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$, which satisfies Assumption (A3), cannot be identified by group symmetries of the network for any choice of the positive weights $\alpha_1, \alpha_2, \alpha_3, \alpha_4, \beta_1,$ and β_2 . The manifold $\mathcal{S}_{\mathcal{P}}$ is invariant whenever the oscillators' natural frequencies satisfy Assumption (A2). Thus, this example shows that network symmetries are not necessary for cluster synchronization of Kuramoto oscillators.

(A2) Given $\mathcal{P} = \{\mathcal{P}_1, \dots, \mathcal{P}_m\}$, the natural frequencies satisfy $\omega_i = \omega_j$ for every $i, j \in \mathcal{P}_k$ and $k \in \{1, \dots, m\}$.³

(A3) The network weights satisfy $\sum_{k \in \mathcal{P}_\ell} a_{ik} - a_{jk} = 0$ for every $i, j \in \mathcal{P}_z$ and $z, \ell \in \{1, \dots, m\}$, with $z \neq \ell$.

Thus, in the remainder of this paper, we assume that (A2) and (A3) are satisfied for the network partition being considered.

Remark 1. (*Network symmetries, equitable partitions, and balanced weights*): Conditions to ensure the invariance of the cluster synchronization manifold have been linked to network symmetries, which are defined by the group comprising all node permutations that leave the network topology unchanged, e.g., see, [31], [32], and [44]. In Fig. 1, we propose a network with two clusters, which are not defined by any group symmetry, that satisfies Assumption (A3), and thus, admits an invariant cluster synchronization manifold. This example shows that cluster synchronization of Kuramoto oscillators does not require symmetric networks. Our Assumption (A3), and in fact, the equivalent notion of external equitable partition [41], is less restrictive than requiring partitions satisfying group symmetries [50]–[52]. Finally, Assumptions (A2) and (A3) are necessary when the natural frequencies in distinct clusters are sufficiently different (see [42] and Remark 2). ■

Remark 2. (*Invariance of submanifolds of $\mathcal{S}_{\mathcal{P}}$*): When the network of oscillators is cluster-synchronized (i.e., $\theta(t) \in \mathcal{S}_{\mathcal{P}}$ for all $t \geq 0$), submanifolds of $\mathcal{S}_{\mathcal{P}}$ may appear whenever the phases belonging to two (or more) disjoint clusters have equal values (see Fig. 2). Interestingly, the example in Fig. 2 also points out that Assumption (A3) may not be necessary for the invariance of $\mathcal{S}_{\mathcal{P}}$ if the clusters do not evolve with different frequencies (see [42, Assumption (A1)]). In what follows, we show that, if the natural frequencies of the oscillators in disjoint clusters are sufficiently different, then invariant, and hence, stable, submanifolds cannot exist. To see this, assume that the phases of the disjoint clusters \mathcal{P}_ℓ and \mathcal{P}_z remain equal over time. Then, using Assumptions (A2) and (A3), the dynamics

³This condition is necessary for $\mathcal{S}_{\mathcal{P}}$ to be forward invariant, and thus, stable [42], and is motivated by observed synchronization phenomena, e.g., see [49].

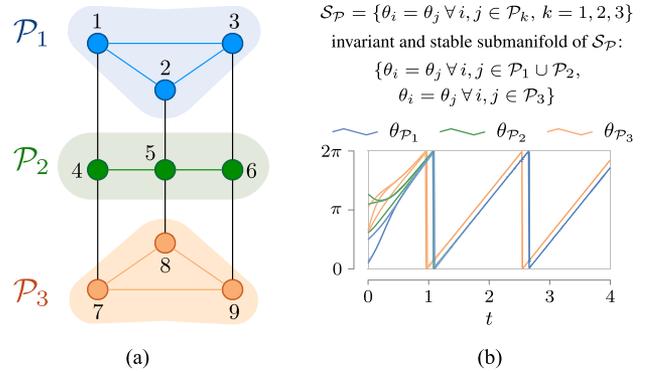


Fig. 2. (a) Network with partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3\}$. As shown in (b), the phases of the oscillators in \mathcal{P}_1 and \mathcal{P}_2 have the same value over time, showing that a submanifold of $\mathcal{S}_{\mathcal{P}}$ is invariant and stable. For this simulation, we use $\omega_1 = 4, \omega_2 = 2, \omega_3 = 6, a_{14} = 3,$ and $a_{47} = 5$.

$$\begin{aligned} \dot{\theta}_\ell - \dot{\theta}_z &= \omega_\ell - \omega_z + \sum_{k=1}^m \left[\left(\sum_{r \in \mathcal{P}_k} a_{\ell r} \right) \sin(\theta_k - \theta_\ell) \right. \\ &\quad \left. - \left(\sum_{r \in \mathcal{P}_k} a_{zr} \right) \sin(\theta_k - \theta_z) \right] \end{aligned} \quad (2)$$

must be identically zero, where θ_i denotes the phase of any oscillator in \mathcal{P}_i . Clearly, if the following inequality holds:

$$|\omega_\ell - \omega_z| > 2(m-2) \max_{k \neq \ell, z} \left\{ \sum_{r \in \mathcal{P}_k} a_{\ell r}, \sum_{r \in \mathcal{P}_k} a_{zr} \right\} \quad (3)$$

Equation (2) cannot vanish, and consequently, the clusters \mathcal{P}_ℓ and \mathcal{P}_z cannot evolve with the same phases when the network is cluster synchronized.⁴ More generally, if condition (3) is satisfied for all pairs of clusters, then invariant, and hence, stable, cluster synchronization submanifolds cannot exist. ■

We conclude with an example showing that the synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is, in general, not globally asymptotically stable due to the existence of multiple invariant sets.

Example 1. (*Multiple invariant sets*): Consider a Kuramoto network with $2N$ oscillators ($N \geq 2$) and with an adjacency matrix defined as follows⁵ (see Fig. 3(a) for the case $N = 5$):

$$a_{ij} = \begin{cases} 1, & \text{if } |i - j| \leq 2 \\ 0, & \text{otherwise} \end{cases}$$

with $i, j \in \{1, \dots, 2N\}$ (and the convention $2N + \ell \triangleq \ell, -\ell \triangleq 2N + \ell - 1$, for $\ell \in \{1, 2\}$). Let $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$, with $\mathcal{P}_1 = \{1, 3, \dots, 2N - 1\}, \mathcal{P}_2 = \{2, 4, \dots, 2N\}$, and define

$$\mathcal{M}_{\mathcal{P}} = \{\theta \in \mathbb{T}^{2N} : \theta_{i+2} = \theta_i + 2\pi/N, i = 1, \dots, 2N - 2\}.$$

It can be verified that Assumption (A3) is satisfied, and that the set $\mathcal{S}_{\mathcal{P}}$ is invariant whenever the natural frequencies satisfy Assumption (A2). Yet, the set $\mathcal{S}_{\mathcal{P}}$ is not the only invariant set. In fact, $\mathcal{M}_{\mathcal{P}}$ is also invariant (we prove this by showing that

⁴In (3), we have $(m-2)$ because for $k = z, \ell$, the sine terms in (2) vanish.

⁵This analysis extends directly to arbitrary weights $a_{ij} = a, a \in \mathbb{R}_{>0}$.

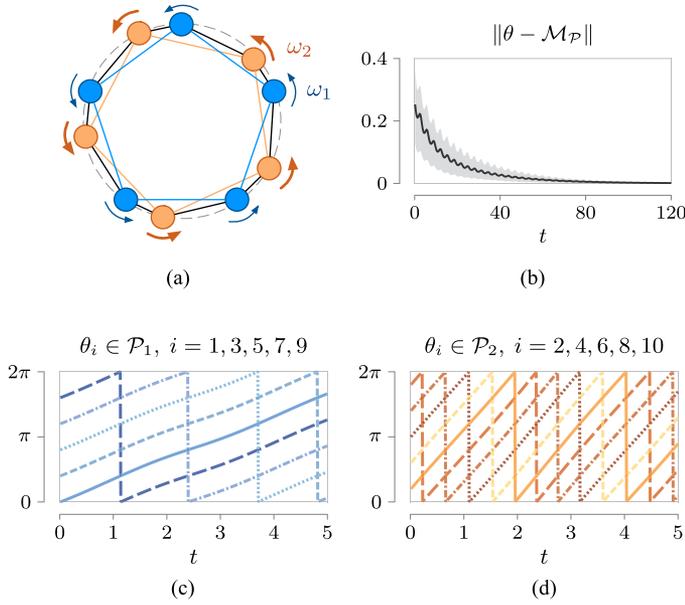


Fig. 3. (a) Network in Example 1 for the case $N = 5$. The nodes belonging to partition \mathcal{P}_1 are blue and have natural frequency $\omega_1 = 1$, while the nodes belonging to partition \mathcal{P}_2 are orange and have $\omega_2 = 3$. (b) Stability of the set $\mathcal{M}_{\mathcal{P}}$ via numerical simulations. We performed 10^3 iterations, each one with $\theta(0)$ chosen randomly within an angle of ± 0.01 [rad] from $\mathcal{M}_{\mathcal{P}}$. The thick line represents the mean value among all simulations of the 2-norm distance between θ and $\mathcal{M}_{\mathcal{P}}$, while the faded area represents the smallest and largest value of the 2-norm distance between θ and $\mathcal{M}_{\mathcal{P}}$. (c) and (d) Invariance of the set $\mathcal{M}_{\mathcal{P}}$ as the phases in the clusters \mathcal{P}_1 and \mathcal{P}_2 evolve, respectively, with the same frequencies.

$$\dot{\theta}_i = \dot{\theta}_{i+2} \text{ when } \theta_i, \theta_{i+2} \in \mathcal{M}_{\mathcal{P}}$$

$$\begin{aligned} \dot{\theta}_i &= \omega_i + \sin(\theta_{i-2} - \theta_i) + \sin(\theta_{i+2} - \theta_i) \\ &\quad + \sin(\theta_{i-1} - \theta_i) + \sin(\theta_{i+1} - \theta_i) \\ &= \omega_i + \sin(\theta_i - \theta_{i+2}) + \sin(\theta_{i+4} - \theta_{i+2}) \\ &\quad + \sin(\theta_{i+1} - \theta_{i+2}) + \sin(\theta_{i+3} - \theta_{i+2}) = \dot{\theta}_{i+2} \end{aligned}$$

where we have used the fact that $\theta_{i+2} - \theta_i = 2\pi/N$, and $\omega_i = \omega_j$ for all i, j in the same cluster. Further, it can be verified numerically that, depending on the number of oscillators N , the set $\mathcal{M}_{\mathcal{P}}$ is also locally stable [see Fig. 3(b)]. We conclude that the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is not, in general, globally asymptotically stable. In what follows, we derive conditions guaranteeing local stability of $\mathcal{S}_{\mathcal{P}}$. ■

III. CONDITIONS FOR THE STABILITY OF THE CLUSTER SYNCHRONIZATION MANIFOLD

In this section, we derive sufficient conditions for the local exponential stability of the cluster synchronization manifold. Define the phase difference $x_{ij} = \theta_j - \theta_i$, and notice that

$$\dot{x}_{ij} = \omega_j - \omega_i + \sum_{z=1}^n [a_{jz} \sin(x_{jz}) - a_{iz} \sin(x_{iz})]. \quad (4)$$

Given a partition $\mathcal{P} = \{\mathcal{P}_1, \dots, \mathcal{P}_m\}$ of the set \mathcal{V} in the graph \mathcal{G} , we define the following graphs (see also Example 2):

- 1) the graph of the k th cluster, with $k \in \{1, \dots, m\}$, $\mathcal{G}_k = (\mathcal{P}_k, \mathcal{E}_k)$, where $\mathcal{E}_k = \{(i, j) : (i, j) \in \mathcal{E}, i, j \in \mathcal{P}_k\}$;
- 2) a spanning tree $\mathcal{T}_k = (\mathcal{P}_k, \mathcal{E}_{\text{span},k})$ of \mathcal{G}_k ;⁶
- 3) a spanning tree $\mathcal{T} = (\mathcal{V}, \mathcal{E}_{\mathcal{T}})$ of \mathcal{G} with $\mathcal{E}_{\mathcal{T}} = \bigcup_{k=1}^m \mathcal{E}_{\text{span},k} \cup \mathcal{E}_{\text{inter}}$, where $\mathcal{E}_{\text{inter}}$ satisfies $|\mathcal{E}_{\text{inter}}| = m - 1$.

Further, we define the following vectors of phase differences:

- 4) $x_{\text{intra}}^{(k)} = [x_{ij}]$ for all $(i, j) \in \mathcal{E}_{\text{span},k}$ with $i < j$;
- 5) $x_{\text{intra}} = [x_{\text{intra}}^{(1)\top}, \dots, x_{\text{intra}}^{(m)\top}]^{\top}$;
- 6) $x_{\text{inter}} = [x_{ij}]$ for all $(i, j) \in \mathcal{E}_{\text{inter}}$ with $i < j$.

It should be noticed that the vectors $x_{\text{intra}}^{(k)}$, x_{intra} , and x_{inter} contain, respectively, $n_{\text{intra},k} = |\mathcal{P}_k| - 1$, $n_{\text{intra}} = n - m$, and $n_{\text{inter}} = m - 1$ entries. Notice that every phase difference can be computed as a linear function of x_{intra} and x_{inter} . To see this, let $i, j \in \mathcal{V}$, and let $p(i, j) = \{p_1, \dots, p_\ell\}$ be the unique path on \mathcal{T} from i to j . Define $\text{diff}(p(i, j)) = \sum_{k=1}^{\ell-1} s_k$, where $s_k = x_{p_k p_{k+1}}$ if $p_k < p_{k+1}$, and $s_k = -x_{p_{k+1} p_k}$ otherwise. Then, $x_{ij} = \text{diff}(p(i, j))$, and the vectors x_{intra} and x_{inter} contain a smallest set of phase differences that can be used to quantify the synchronization among all of the oscillators in the network.

Let $B = [b_{k\ell}] \in \mathbb{R}^{|\mathcal{V}| \times |\mathcal{E}|}$ denote the oriented incidence matrix of the graph $\mathcal{G} = (\mathcal{V}, \mathcal{E})$, where ℓ corresponds to the edge $(i, j) \in \mathcal{E}$, $b_{k\ell} = 1$ if node k is the sink of the edge ℓ , $b_{k\ell} = -1$ if k is the source of ℓ , and $b_{k\ell} = 0$ otherwise.⁷ Further, let B_k and $B_{\text{span},k}$ denote the incidence matrices of \mathcal{G}_k and \mathcal{T}_k , respectively. Notice that $B_{\text{span},k}$ is of full rank because it is the incidence matrix of an acyclic graph (tree) [53, Th. 8.3.1]. Let $T_{\text{intra},k} = B_k^{\top} (B_{\text{span},k}^{\top})^{\dagger}$ be the unique matrix that maps the phase differences contained in $x_{\text{intra}}^{(k)}$ to all intracluster phase differences in the k th cluster. That is,

$$x^{(k)} = T_{\text{intra},k} x_{\text{intra}}^{(k)} \quad (5)$$

where $x^{(k)}$ contains all phase differences in the cluster \mathcal{P}_k .

We conclude this part by rewriting the intracluster dynamics in a form that will be useful to prove our results. In particular, from the aforementioned discussion and for an intracluster phase difference x_{ij} of $x_{\text{intra}}^{(k)}$, we rewrite (4) as

$$\begin{aligned} \dot{x}_{ij} &= \underbrace{\sum_{z \in \mathcal{P}_k} [a_{jz} \sin(\text{diff}(p(j, z))) - a_{iz} \sin(\text{diff}(p(i, z)))]}_{F_{ij}^{(k)}(x_{\text{intra}}^{(k)})} \\ &\quad + \underbrace{\sum_{z \notin \mathcal{P}_k} [a_{jz} \sin(\text{diff}(p(j, z))) - a_{iz} \sin(\text{diff}(p(i, z)))]}_{G_{ij}^{(k)}(x_{\text{intra}}, x_{\text{inter}})} \end{aligned} \quad (6)$$

which leads to

$$\dot{x}_{\text{intra}}^{(k)} = F^{(k)}(x_{\text{intra}}^{(k)}) + G^{(k)}(x_{\text{intra}}, x_{\text{inter}}) \quad (7)$$

⁶We assume that \mathcal{G} and its subgraphs \mathcal{G}_k are connected. This guarantees the existence of the (connected) spanning trees defined in 2) and 3). A graph is connected if there exists a path between any pair of nodes [53].

⁷Node i is the source (resp., sink) of (i, j) if $i < j$ (resp., $i > j$).

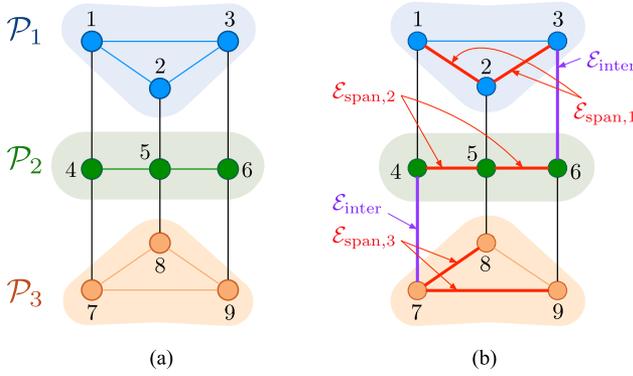


Fig. 4. Graph-theoretic definitions introduced in Section III for a network of nine Kuramoto oscillators (see also Example 2). (a) Partitions $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3\}$. (b) $\mathcal{E}_{\text{span},1}$, $\mathcal{E}_{\text{span},2}$, and $\mathcal{E}_{\text{span},3}$ represent (in red) the edges of the intracluster spanning trees T_1 , T_2 , and T_3 , respectively, while the edges belonging to the set $\mathcal{E}_{\text{inter}}$ are depicted in purple.

where $F^{(k)}$ is the vector of $F_{ij}^{(k)}$ and $G^{(k)}$ is the vector of $G_{ij}^{(k)}$, for all $(i, j) \in \mathcal{E}_{\text{span},k}$ with $i < j$. Finally, by concatenating the dynamics (7) for all clusters, we obtain

$$\dot{x}_{\text{intra}} = F(x_{\text{intra}}) + G(x_{\text{intra}}, x_{\text{inter}}). \quad (8)$$

Example 2. (Illustration of the definitions): We provide here an illustrative example of the definitions introduced in this section. Consider the network in Fig. 4(a) with partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3\}$, where $\mathcal{P}_1 = \{1, 2, 3\}$, $\mathcal{P}_2 = \{4, 5, 6\}$, and $\mathcal{P}_3 = \{7, 8, 9\}$. Fig. 4(b) illustrates the definitions of spanning trees, together with the edge sets $\mathcal{E}_{\text{span},k}$ ($k = 1, 2, 3$), and the intercluster edges in $\mathcal{E}_{\text{inter}} = \{(3, 6), (4, 7)\}$. The vectors of intracluster differences read as $x_{\text{intra}}^{(1)} = [x_{12} \ x_{23}]^T$, $x_{\text{intra}}^{(2)} = [x_{45} \ x_{56}]^T$, and $x_{\text{intra}}^{(3)} = [x_{78} \ x_{79}]^T$, whereas the vector of intercluster differences reads as $x_{\text{inter}} = [x_{36} \ x_{47}]^T$.

For the partition \mathcal{P}_1 , order the edges as $\ell_1 = (1, 2)$, $\ell_2 = (1, 3)$, and $\ell_3 = (2, 3)$. Then, a spanning tree is $T_1 = (\mathcal{P}_1, \mathcal{E}_{\text{span},1})$, with $\mathcal{E}_{\text{span},1} = \{(1, 2), (2, 3)\}$, and the (oriented) incidence matrices B_1 of \mathcal{G}_1 and $B_{\text{span},1}$ of T_1 are

$$B_1 = \begin{bmatrix} -1 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 1 \end{bmatrix}, \quad B_{\text{span},1} = \begin{bmatrix} -1 & 0 \\ 1 & -1 \\ 0 & 1 \end{bmatrix}.$$

Finally, the matrix $T_{\text{intra},1} = B_1^T (B_{\text{span},1}^T)^{\dagger}$ satisfies

$$T_{\text{intra},1} = \begin{bmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \end{bmatrix}^T.$$

A. Asymptotic Stability of $\mathcal{S}_{\mathcal{P}}$ Via the Perturbation Theory

In what follows, we will make use of the perturbation theory of dynamical systems to provide our first stability condition. We first introduce the following instrumental result.

Lemma 3.1. (Properties of intracluster dynamics): The intracluster dynamics (8) satisfies the following properties.

- 1) The Jacobian matrix J_{intra} of $F(x_{\text{intra}})$ computed at the origin is Hurwitz stable and can be written as

$$J_{\text{intra}} = \left. \frac{\partial F(x_{\text{intra}})}{\partial x_{\text{intra}}} \right|_{x_{\text{intra}}=0} = \text{blk-diag}(J_1, \dots, J_m) \quad (9)$$

where for $k \in \{1, \dots, m\}$, $T_{\text{intra},k}$ is as in (5) and

$$J_k = -B_{\text{span},k}^T B_k \text{diag}(\{a_{ij}\}_{(i,j) \in \mathcal{E}_k}) T_{\text{intra},k}. \quad (10)$$

Thus, the origin is an exponentially stable equilibrium of the system $\dot{x}_{\text{intra}} = F(x_{\text{intra}})$.

- 2) There exist constants $\gamma^{(k\ell)} \in \mathbb{R}_{>0}$ such that

$$\|G^{(k)}(x_{\text{intra}}, x_{\text{inter}})\| \leq \sum_{\ell=1}^m \gamma^{(k\ell)} \|x_{\text{intra}}^{(\ell)}\| \quad (11)$$

for all $k, \ell \in \{1, \dots, m\}$. Specifically

$$\gamma^{(k\ell)} = 2 \max_r n_{\text{intra},r} \tilde{\gamma}^{(k\ell)} \quad (12)$$

where for any $i \in \mathcal{P}_k$

$$\tilde{\gamma}^{(k\ell)} = \begin{cases} \sum_{j \in \mathcal{P}_\ell} a_{ij}, & \text{if } \ell \neq k \\ \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{j \in \mathcal{P}_\ell} a_{ij}, & \text{otherwise.} \end{cases} \quad (13)$$

As formalized in the next theorem, Lemma 3.1, together with results on the stability of perturbed systems [54, ch. 9], implies that the origin of (8), and thus, the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$, is exponentially stable for some choices of the network weights. Recall that an M -matrix is a real nonsingular matrix $A = [a_{ij}]$ such that $a_{ij} \leq 0$ for all $i \neq j$ and all leading principal minors are positive [55, ch. 2.5].

Theorem 3.2. (Sufficient condition on network weights for the stability of $\mathcal{S}_{\mathcal{P}}$): Let $\mathcal{S}_{\mathcal{P}}$ be the cluster synchronization manifold associated with a partition $\mathcal{P} = \{\mathcal{P}_1, \dots, \mathcal{P}_m\}$ of the network \mathcal{G} of Kuramoto oscillators. Let $\gamma^{(k\ell)}$ be the constants defined in (12). Define the matrix $S \in \mathbb{R}^{m \times m}$ as

$$S = [s_{k\ell}] = \begin{cases} \lambda_{\max}^{-1}(X_k) - \gamma^{(kk)} & \text{if } k = \ell \\ -\gamma^{(k\ell)} & \text{if } k \neq \ell \end{cases} \quad (14)$$

where $X_k > 0$ is such that $J_k^T X_k + X_k J_k = -I$, with J_k as in (10). If S is an M -matrix, then the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is locally exponentially stable.

Remark 3. (Family of bounds): In (14), the matrices X_k can be selected as the solutions to the Lyapunov equations $J_k^T X_k + X_k J_k = -Q_k$, for arbitrary positive definite matrices Q_k . Yet, selecting $Q_k = I$ for all k yields a tighter stability bound. This follows because 1) if S is an M -matrix, then $S + \Delta$ remains an M -matrix whenever Δ is a nonnegative diagonal matrix [55, Th. 2.5.3], and 2) the ratio $\lambda_{\min}(Q_k)/\lambda_{\max}(X_k)$ is maximal whenever $Q_k = I$ [54, Exercise 9.1].

Theorem 3.2 describes a sufficient condition on the network weights for the stability of the cluster synchronization manifold. Loosely speaking, the cluster synchronization manifold is exponentially stable when the intracluster coupling (measured by $\lambda_{\max}^{-1}(X_k) - \gamma^{(kk)}$) is sufficiently stronger than the perturbation induced by the intercluster connections (measured by $\gamma^{(k\ell)}$). In

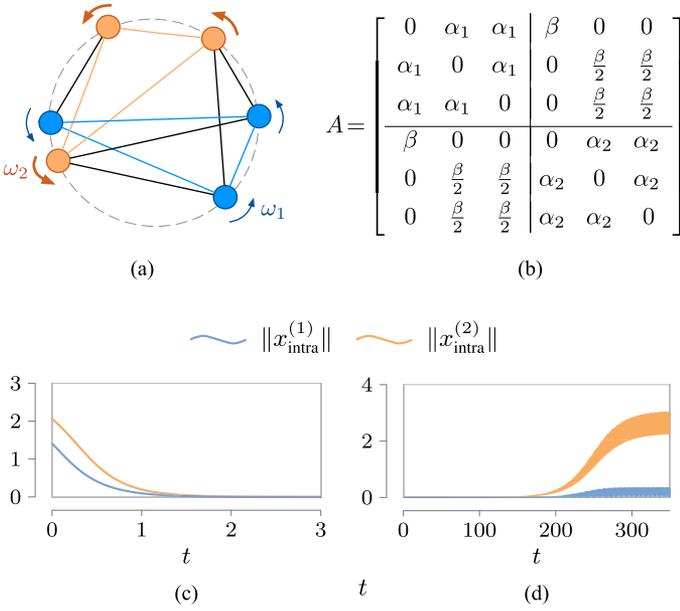


Fig. 5. (a) Network of six Kuramoto oscillators in Example 3. We identify the clusters \mathcal{P}_1 and \mathcal{P}_2 in blue and orange, respectively. (b) Adjacency matrix of the network in (a). The parameters α_1, α_2 , and β represent the intracluster and intercluster weights, respectively. (c) Stability of the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ for $\alpha_1 = \alpha_2 = 1$ and $\beta = 0.1$, as predicted by Theorem 3.2. (d) $\mathcal{S}_{\mathcal{P}}$ is unstable when $\alpha_1 = \beta = 1$ and $\alpha_2 = 0.001$.

particular, the term $\lambda_{\max}^{-1}(X_k)$ is proportional to the intracluster weights and it is implicitly related to the network topology. In fact, the matrix X_k is the solution of a Lyapunov's equation containing J_k , whose spectrum coincides with the stable eigenvalues of the negative Laplacian matrix of the k th cluster. We refer the interested reader to the proof of Lemma 3.1. Finally, we remark that a result akin to Theorem 3.2 has been derived in [56], although for interconnected systems whose coupling functions are required to satisfy certain assumptions that fail to hold in the Kuramoto model.

Example 3. (*Tradeoff between intra- and intercluster weights*): Consider the network in Fig. 5(a) with partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$, where $\mathcal{P}_1 = \{1, 2, 3\}$ and $\mathcal{P}_2 = \{4, 5, 6\}$, natural frequencies $\omega_1 = 1$ and $\omega_2 = 6$ for the oscillators in \mathcal{P}_1 and \mathcal{P}_2 , and adjacency matrix as in Fig. 5(b). The parameters $\alpha_1, \alpha_2 \in \mathbb{R}_{>0}$, and $\beta \in \mathbb{R}_{>0}$ denote the strength of the intra- and intercluster coupling, respectively. Let $\alpha_1 = \alpha_2$, and construct the matrix S as in Theorem 3.2 as

$$S = \begin{bmatrix} \lambda_{\max}^{-1}(X_1) - \gamma_{11} & -\gamma_{12} \\ -\gamma_{12} & \lambda_{\max}^{-1}(X_2) - \gamma_{22} \end{bmatrix}$$

where $X_k \succ 0$ is such that $J_k^T X_k + X_k J_k = -I$, $\lambda_{\max}^{-1}(X_1) = \lambda_{\max}^{-1}(X_2) = 2\alpha_1$, and from (12), $\gamma_{ij} = 4\beta$ for all i, j . By inspecting all leading principal minors, S is an M -matrix if $\alpha_1/\beta > 4$, and the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is exponentially stable [see Fig. 5(c)]. We remark that, when $\alpha_1 \neq \alpha_2$, the synchronization manifold $\mathcal{S}_{\mathcal{P}}$ can become unstable, as we verify numerically in Fig. 5(d). ■

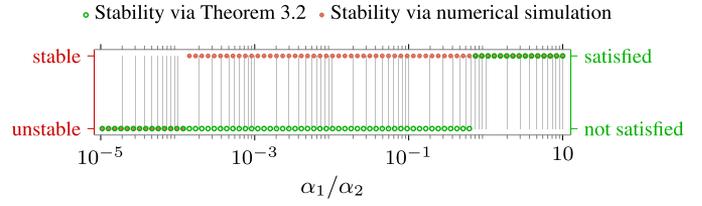


Fig. 6. The condition in Theorem 3.2 leads to conservative stability bounds. For the network in Example 3, we let $\beta = 0.1$ and plot, as a function of the ratio α_1/α_2 , the stable configurations predicted by Theorem 3.2 (green) and those found numerically. For each value of α_1/α_2 , we assess numerical stability by making use of the Floquet stability theory [57, ch. 5] and by resorting to statement 1) in Lemma 3.4. This is possible because the partition in Example 3 has only two clusters.

The stability condition in Theorem 3.2 depends only on the network weights, and typically leads to conservative bounds (see also Fig. 6). To derive refined stability conditions, we next characterize how the natural frequencies of the oscillators affect stability of the cluster synchronization manifold.

B. Asymptotic Stability of $\mathcal{S}_{\mathcal{P}}$ When the Oscillators' Natural Frequencies are Sufficiently Different

Natural frequencies play a fundamental role for full and cluster synchronization of Kuramoto oscillators. However, while heterogeneity of the natural frequencies typically impedes full synchronization [47], we will show that cluster synchronization is in fact facilitated when the oscillators in different clusters have sufficiently different natural frequencies. We start with an asymptotic result that is valid for arbitrary networks and partitions, and then, improve our results for the case of partitions containing only two clusters.

Theorem 3.3. (*Stability of $\mathcal{S}_{\mathcal{P}}$ for large natural frequency differences*): Let $\mathcal{S}_{\mathcal{P}}$ be the cluster synchronization manifold associated with a partition $\mathcal{P} = \{\mathcal{P}_1, \dots, \mathcal{P}_m\}$ of the network \mathcal{G} of Kuramoto oscillators. Let $\omega_i \in \mathbb{R}_{>0}$ be the natural frequency of the oscillators in the cluster \mathcal{P}_i , with $i \in \{1, \dots, m\}$. In the limit $|\omega_i - \omega_j| \rightarrow \infty$, for all $i, j \in \{1, \dots, m\}, i \neq j$, the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is locally exponentially stable.

Theorem 3.3 shows that heterogeneity of the natural frequencies of the oscillators in different clusters facilitates cluster synchronization, independently of the network weights. We remark that a similar behavior was also identified in [58], albeit with a different method and definition of synchronization.

We next improve upon Theorem 3.3 by analyzing the case where the natural frequencies are finite and the partition \mathcal{P} contains only two clusters. To this aim, let $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$ and assume, without loss of generality, that $\omega_2 \geq \omega_1$, where ω_i is the natural frequency of the oscillators in \mathcal{P}_i . Define

$$\bar{\omega} = \omega_2 - \omega_1, \text{ and } \bar{a} = \sum_{k \in \mathcal{P}_2} a_{ik} + \sum_{k \in \mathcal{P}_1} a_{jk}$$

for any $i \in \mathcal{P}_1$ and $j \in \mathcal{P}_2$. The next result characterizes the intercluster phase difference when the network evolves on the cluster synchronization manifold.

Lemma 3.4. (*Nominal intercluster difference*): Let $\mathcal{S}_{\mathcal{P}}$ be the cluster synchronization manifold associated with a partition

$\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$ of the network \mathcal{G} of Kuramoto oscillators. Let $\theta(0) \in \mathcal{S}_{\mathcal{P}}$ (equivalently, $x_{\text{intra}}(0) = 0$). Then, if $x_{\text{intra}} = 0$ at all times and $\bar{\omega} > \bar{a}$

$$x_{\text{inter}}(t) = \begin{cases} h(t), & \text{if } t \neq t_0 + kT, k \in \mathbb{Z} \\ \pi, & \text{if } t = t_0 + kT, k \in \mathbb{Z} \end{cases} \triangleq x_{\text{nom}}(t) \quad (15)$$

where

$$h(t) = 2 \tan^{-1} \left(\frac{\bar{a} + \sqrt{\bar{\omega}^2 - \bar{a}^2} \tan\left(\frac{\sqrt{\bar{\omega}^2 - \bar{a}^2}}{2}(t + \tau)\right)}{\bar{\omega}} \right)$$

$t_0 = -\tau + \pi/\sqrt{\bar{\omega}^2 - \bar{a}^2}$, $T = 2\pi/\sqrt{\bar{\omega}^2 - \bar{a}^2}$, and $\tau \in \mathbb{R}$ is a constant that depends only on $\theta(0)$. Moreover

- 1) x_{nom} is T -periodic with zero time average;
- 2) the following inequality holds:

$$\left| \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau \right| \leq \frac{1}{\bar{a}} \log \left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right). \quad (16)$$

Remark 4. (Constant versus time-varying intercluster difference): The values of $\bar{\omega}$ and \bar{a} determine the behavior of the intercluster phase difference. In particular, if $\bar{\omega} < \bar{a}$, then the intercluster difference evolves as in (15).⁸ If $\bar{\omega} = \bar{a}$, (4) reduces to $\dot{x}_{\text{inter}} = \bar{a} - \bar{a} \sin(x_{\text{inter}})$, which can be integrated as follows:

$$\begin{aligned} \bar{a}t &= \int_{x_{\text{inter}}(0)}^{x_{\text{inter}}(t)} (1 - \sin(s))^{-1} ds \\ \bar{a}t &= \frac{2 \sin(x_{\text{inter}}(t)/2)}{\cos(x_{\text{inter}}(t)/2) - \sin(x_{\text{inter}}(t)/2)} + \tau. \end{aligned} \quad (17)$$

By substitution, it can be verified that

$$x_{\text{inter}}(t) = 2 \cos^{-1} \left(\frac{\bar{a}t - \tau + 2}{\sqrt{2(\bar{a}t - \tau + 1)^2 + 2}} \right)$$

satisfies (17). In both cases ($\bar{\omega} \leq \bar{a}$), x_{inter} converges to the constant value $2 \tan^{-1}((\bar{a} - \sqrt{\bar{a}^2 - \bar{\omega}^2})/\bar{\omega})$ as t increases to infinity. In other words, if $\bar{\omega} \leq \bar{a}$, then the phases of the oscillators in the two clusters evolve with the same frequency, and the oscillators are phase locked (see Fig. 7(a) and [47, Remark 1]). Instead, if $\bar{\omega} > \bar{a}$, the clusters evolve with different frequencies, and the intercluster phase difference follows a limit cycle (see Fig. 7(b) and [54, ch. 2]). ■

In the remainder of this section, we assume that $\bar{\omega} > \bar{a}$ so that the clusters evolve with different frequencies (see Remark 4). Leveraging Lemma 3.4, we next present a refined condition for the stability of the cluster synchronization manifold.

Theorem 3.5. (Sufficient condition on network weights and natural frequencies for the stability of $\mathcal{S}_{\mathcal{P}}$): Let $\mathcal{S}_{\mathcal{P}}$ be the cluster synchronization manifold associated with a partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$ of the network \mathcal{G} of Kuramoto oscillators. Let $\omega_i \in \mathbb{R}_{>0}$ be the natural frequency of the oscillators in the cluster \mathcal{P}_i , with $i \in \{1, 2\}$. Let J_{intra} be as in Lemma 3.1, and

⁸In fact, $\sqrt{\bar{\omega}^2 - \bar{a}^2}$ becomes a complex number, and by recalling that $\tan(i\alpha) = i \tanh(\alpha)$, where $\alpha \in \mathbb{R}$, in (15), we have $x_{\text{inter}}(t) = 2 \tan^{-1}((\bar{a} - \sqrt{\bar{a}^2 - \bar{\omega}^2} \tanh(\sqrt{\bar{a}^2 - \bar{\omega}^2}(t + \tau)/2))/\bar{\omega})$.

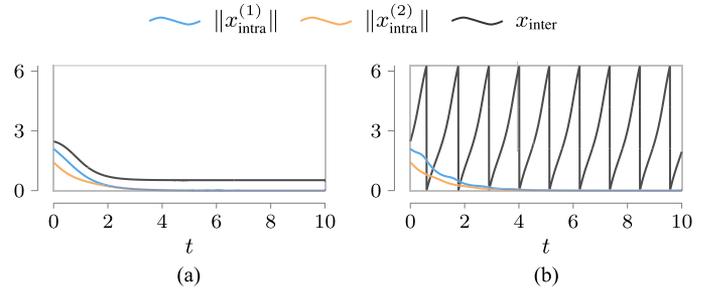


Fig. 7. For the network in Example 3 with $\alpha_1 = \alpha_2 = \beta = 1$, $\bar{a} = 2$ and $\bar{\omega} = 1$, (a) shows that the clusters are synchronized (as $\|x_{\text{intra}}^{(1)}\|$ and $\|x_{\text{intra}}^{(2)}\|$ converge to zero), yet all oscillators remain phase locked (x_{inter} converges to a constant). Instead, (b) shows that the intercluster difference follows a limit cycle when $\alpha_1 = \alpha_2 = \beta = 1$, $\bar{a} = 2$, and $\bar{\omega} = 6$.

$J_{\text{inter}} = \partial G(x_{\text{intra}}, x_{\text{inter}})/\partial x_{\text{intra}}$ along the trajectory $x_{\text{intra}} = 0$ and $x_{\text{inter}} = x_{\text{nom}}$. The cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is locally exponentially stable if the following inequality holds:

$$\left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right)^{\frac{2}{\bar{a}} \|J_{\text{inter}}\|} < 1 + \frac{1}{2\lambda_{\max}(X) \|J_{\text{intra}}\|} \quad (18)$$

where $X \succ 0$ is the solution of $J_{\text{intra}}^T X + X J_{\text{intra}} = -I$.

Theorem 3.5 provides a quantitative condition on the network weights and the natural frequencies of the oscillators to ensure the stability of the cluster synchronization manifold. It can be shown that 1) when the intercluster weights decrease to zero ($\bar{a} \rightarrow 0$) and $\bar{\omega}$ remains bounded, then $\|J_{\text{inter}}\|/\bar{a}$ remains bounded, the left-hand side of (18) converges to 1, and the inequality is automatically satisfied, and 2) when $\bar{\omega}$ grows ($\bar{\omega} \rightarrow \infty$) and the intercluster weights remain bounded, the left-hand side of (18) converges to 1 and the inequality is automatically satisfied. The role of the intracluster connections on the stability of $\mathcal{S}_{\mathcal{P}}$ cannot be evaluated directly from (18) because of the dependence of the right-hand side on $\lambda_{\max}(X)$. The following result, however, suggests that the synchronization manifold may remain exponentially stable when the intracluster weights are homogeneous, independently of the intercluster weights and the natural frequencies.

Theorem 3.6. (Stability of $\mathcal{S}_{\mathcal{P}}$ with homogeneous clusters):

Let $\mathcal{S}_{\mathcal{P}}$ be the cluster synchronization manifold associated with a partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2\}$ of the network \mathcal{G} of Kuramoto oscillators. Let $\omega_i \in \mathbb{R}_{>0}$ be the natural frequency of the oscillators in the cluster \mathcal{P}_i , with $i \in \{1, 2\}$. If $J_{\text{intra}} = \alpha I$, for some constant $\alpha \in \mathbb{R}_{<0}$, then the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is locally exponentially stable.

We provide an example that illustrates the stability conditions derived in Theorem 3.5.

Example 4. (Heterogeneity of natural frequencies improves stability of the cluster synchronization manifold): Consider the network of Kuramoto oscillators in Example 3. Fig. 8(a) illustrates that the cluster synchronization manifold is asymptotically stable when the condition in Theorem 3.5 is satisfied. Fig. 8(b) illustrates the tradeoff in the latter stability condition between the natural frequency $\bar{\omega}$ and the intercluster strength measured by β^* , which denotes the largest intercluster weight β (see Ex-

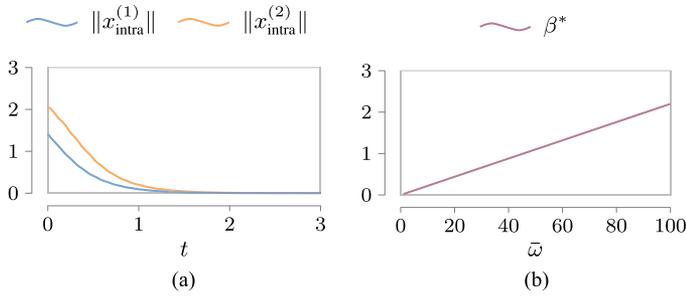


Fig. 8. For the network in Example 3, (a) illustrates the stability of $\mathcal{S}_{\mathcal{P}}$ when $\alpha_1 = \alpha_2 = \beta = \omega_1 = 1$ and $\omega_2 = 47$, as predicted by the condition in Theorem 3.5. For the same network and weights, (b) shows the largest value of intercluster weights β^* that satisfies (18) with equality. As predicted by Theorem 3.3 and Theorem 3.5, stability of the cluster synchronization manifold $\mathcal{S}_{\mathcal{P}}$ is preserved when $\bar{\omega}$ grows with the intercluster weights.

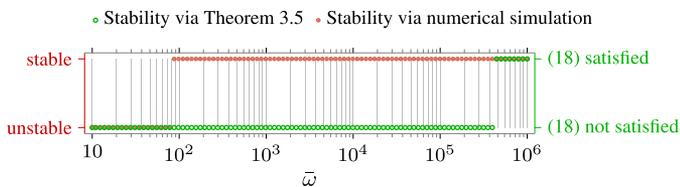


Fig. 9. For the network in Example 3, we let $\alpha_1 = \beta = 1$ and $\alpha_2 = 10^{-4}$ and plot, as a function of $\bar{\omega}$, the stable configurations predicted by Theorem 3.5 (green) and those found numerically. For each value of $\bar{\omega}$, we assess numerical stability (in red) by making use of the Floquet stability theory [57, ch. 5] and by resorting to statement 1) in Lemma 3.4. This is possible because the partition in Example 3 contains two clusters. Although condition (18) is conservative, it captures the effect of large $\bar{\omega}$ on the stability of $\mathcal{S}_{\mathcal{P}}$.

ample 3) such that (18) is still satisfied. Further, we show in Fig. 9 that, while being conservative, condition (18) captures the fact that stability of the cluster synchronization manifold can be recovered by increasing $\bar{\omega}$. Namely, choosing the same network weights that yield instability as in Fig. 5(d), we show that stability of the cluster synchronization manifold is recovered as the difference in natural frequencies grows. ■

We conclude this section with a discussion of cluster synchronization in asymmetric networks and identical nodes.

Remark 5. (*Extension to networks with asymmetric weights*): Symmetry of the network weights is typically exploited to provide conditions for the stability of the full synchronization manifold in networks of Kuramoto oscillators [47]. We rely on the symmetry assumption (A1) to derive statement 1) in Lemma 3.1, which supports our main theorems. However, these results remain valid for bidirected graphs,⁹ provided that the Jacobian J_{intra} can be proven to be Hurwitz. In other words, Assumption (A1) is used to guarantee the stability of the isolated clusters, and not of the cluster configuration. ■

Remark 6. (*Cluster synchronization in networks of identical oscillators*): This paper focuses on heterogeneous oscillators and leverages mismatches in the natural frequencies and

⁹A bidirected graph is a directed graph where $(i, j) \in \mathcal{E}$ implies $(j, i) \in \mathcal{E}$. The adjacency matrix of a bidirected graph needs not be symmetric.

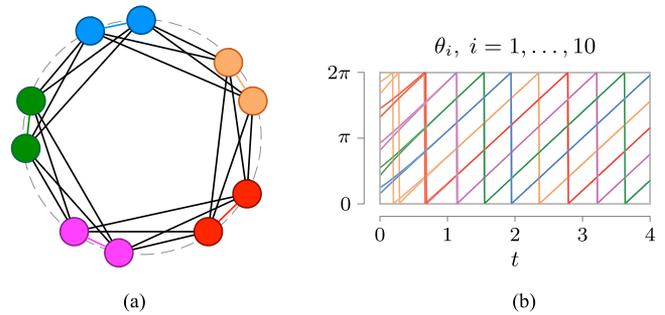


Fig. 10. (a) Network of ten Kuramoto oscillators with partition $\mathcal{P} = \{\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3, \mathcal{P}_4, \mathcal{P}_5\}$, where each cluster is color coded. All oscillators have identical natural frequency $\omega = 3$ and all edges have unit weight. As illustrated in (b), the cluster synchronization manifold associated to \mathcal{P} is stable, showing that cluster synchronization is possible even in networks of identical Kuramoto oscillators with identical edge weights.

the network weights to characterize the stability of the cluster synchronization manifold. Yet, cluster synchronization can also arise in networks of homogeneous Kuramoto oscillators, where all units have equal natural frequencies and all edges have equal weight (e.g., see Fig. 10). With the exception of Theorem 3.3, which is also applicable in the case of identical edge weights, our stability results cannot predict cluster synchronization in networks of identical oscillators, a question that we leave as the subject of future investigation. ■

IV. CONCLUSION AND FUTURE WORK

In this paper, we characterize conditions for the stability of cluster synchronization in networks of oscillators with Kuramoto dynamics, where multiple synchronized groups of oscillators coexist in a connected network. We derive conditions showing that the cluster synchronization manifold is locally exponentially stable when 1) the intracluster coupling is sufficiently stronger than the intercluster coupling, 2) the differences of natural frequencies of the oscillators in disjoint clusters are sufficiently large, or 3) in the case of two clusters, if the intracluster dynamics is homogeneous. To the best of our knowledge, our results are the first to characterize the stability of the cluster synchronization manifold in sparse and weighted networks of heterogeneous Kuramoto oscillators.

Directions of future research include the characterization of tighter stability bounds, the design of methods to control the formation of time-varying synchronized clusters, and the extension of Theorem 3.5 to an arbitrary number of clusters.

APPENDIX

In this section, we provide the proofs of the results presented in Section III, together with some instrumental lemmas.

A. Proofs of the Results in Section III-A

Proof of Lemma 3.1: Proof of statement 1). Notice that the block-diagonal form of the Jacobian matrix J_{intra} follows directly from the form of $F(x_{\text{intra}})$ in (8). Therefore, the stability of J_{intra} is equivalent to the stability of the diagonal blocks J_k . Let $\theta^{(k)}$ be the vector of θ_i , $i \in \mathcal{P}_k$, and by Assumption

(A2), let ω_k be the natural frequency of any oscillator in \mathcal{P}_k . From (1), we write the phase dynamics of the k th cluster as (see [45])

$$\dot{\theta}^{(k)} = \omega_k \mathbf{1} - B_k \text{diag}(\{a_{ij}\}_{(i,j) \in \mathcal{E}_k}) \sin(B_k^\top \theta^{(k)}).$$

Because the phase differences satisfy $x_{\text{intra}}^{(k)} = B_{\text{span},k}^\top \theta^{(k)}$ and $x^{(k)} = B_k^\top \theta^{(k)}$, we have

$$\dot{x}_{\text{intra}}^{(k)} = -B_{\text{span},k}^\top B_k \text{diag}(\{a_{ij}\}_{(i,j) \in \mathcal{E}_k}) \sin(x^{(k)}) \quad (19)$$

where we have used the property $B_{\text{span},k}^\top \mathbf{1} = 0$. Using (5), the Jacobian matrix of (19) computed at $x_{\text{intra}}^{(k)} = 0$ reads as

$$J_k = -B_{\text{span},k}^\top B_k \text{diag}(\{a_{ij}\}_{(i,j) \in \mathcal{E}_k}) T_{\text{intra},k}. \quad (20)$$

Recall that the Laplacian matrix of the graph \mathcal{G}_k satisfies $\mathcal{L}_{\mathcal{G}_k} = B_k \text{diag}(\{a_{ij}\}_{(i,j) \in \mathcal{E}_k}) B_k^\top$, and that, because \mathcal{G}_k is connected, the eigenvalues of $-\mathcal{L}_{\mathcal{G}_k}$ have negative real part, except one single eigenvalue located at the origin with eigenvector $\mathbf{1}$. Define the matrix $W_k = [B_{\text{span},k}^\top \mathbf{1}]^\top$ and notice that, because $B_{\text{span},k}^\top \mathbf{1} = 0$ and $B_{\text{span},k}$ being full column rank [53, Th. 8.3.1], then W_k is invertible and $W_k^{-1} = [(B_{\text{span},k}^\top)^\dagger (\mathbf{1}^\top)^\dagger]$. Therefore, we have

$$W_k (-\mathcal{L}_{\mathcal{G}_k}) W_k^{-1} = \begin{bmatrix} J_k & 0 \\ 0 & 0 \end{bmatrix}$$

where we have used that $T_{\text{intra},k} = B_k^\top (B_{\text{span},k}^\top)^\dagger$ in (20). This shows that J_k contains only the stable eigenvalues of $-\mathcal{L}_{\mathcal{G}_k}$.

Proof of statement 2). Notice that, for any $(j, z) \in \mathcal{E}$ with $j \in \mathcal{P}_k$, $z \in \mathcal{P}_\ell$, and $k \neq \ell$, the difference $\text{diff}(p(j, z))$ in $G_{ij}^{(k)}(x_{\text{intra}}, x_{\text{inter}})$ in (6) can be rewritten as

$$\text{diff}(p(j, z)) = \text{diff}(p(j, k^*)) + \text{diff}(p(k^*, \ell^*)) + \text{diff}(p(\ell^*, z))$$

where k^* and ℓ^* are such that $p(k^*, \ell^*)$ is the shortest path on T connecting the clusters \mathcal{P}_k and \mathcal{P}_ℓ . Then

$$\begin{aligned} G_{ij}^{(k)}(x_{\text{intra}}, x_{\text{inter}}) = & \\ & \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{z \in \mathcal{P}_\ell} [a_{jz} \sin(\text{diff}(p(j, k^*)) + \text{diff}(p(k^*, \ell^*)) + \text{diff}(p(\ell^*, z))) \\ & - a_{iz} \sin(\text{diff}(p(i, k^*)) + \text{diff}(p(k^*, \ell^*)) + \text{diff}(p(\ell^*, z)))]. \end{aligned}$$

Notice that $\text{diff}(p(i, k^*))$ and $\text{diff}(p(j, k^*))$ contain only differences in $x_{\text{intra}}^{(k)}$, and $\text{diff}(p(\ell^*, z))$ only differences in $x_{\text{intra}}^{(\ell)}$.

Notice that $\sin(a + b) = \sin(a) + \delta$, with $|\delta| \leq |b|$.¹⁰ Then

$$\begin{aligned} G_{ij}^{(k)}(x_{\text{intra}}, x_{\text{inter}}) = & \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{z \in \mathcal{P}_\ell} [a_{jz} (\sin(\text{diff}(p(k^*, \ell^*)) + \delta_{jz}) \\ & - a_{iz} (\sin(\text{diff}(p(k^*, \ell^*)) + \delta_{iz})) \\ = & \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \left(\sum_{z \in \mathcal{P}_\ell} [(a_{jz} - a_{iz}) \sin(\text{diff}(p(k^*, \ell^*)))] \right. \\ & \left. + \sum_{z \in \mathcal{P}_\ell} [a_{jz} \delta_{jz} - a_{iz} \delta_{iz}] \right) \stackrel{(A3)}{=} \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{z \in \mathcal{P}_\ell} [a_{jz} \delta_{jz} - a_{iz} \delta_{iz}] \end{aligned}$$

where δ_{jz} and δ_{iz} are upper bounded by $\sqrt{n_{\text{intra},k}} \|x_{\text{intra}}^{(k)}\| + \sqrt{n_{\text{intra},\ell}} \|x_{\text{intra}}^{(\ell)}\|$. Therefore, we have the following bound:

$$\begin{aligned} |G_{ij}^{(k)}| \leq & \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \left(\sum_{z \in \mathcal{P}_\ell} a_{jz} |\delta_{jz}| + \sum_{z \in \mathcal{P}_\ell} a_{iz} |\delta_{iz}| \right) \\ \stackrel{(A3)}{\leq} & 2 \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{z \in \mathcal{P}_\ell} a_{jz} \left(\sqrt{n_{\text{intra},k}} \|x_{\text{intra}}^{(k)}\| + \sqrt{n_{\text{intra},\ell}} \|x_{\text{intra}}^{(\ell)}\| \right) \\ = & 2 \sum_{\ell=1}^m \sqrt{n_{\text{intra},\ell}} \tilde{\gamma}_{ij}^{(k\ell)} \|x_{\text{intra}}^{(\ell)}\| \end{aligned}$$

where

$$\tilde{\gamma}_{ij}^{(k\ell)} = \begin{cases} \sum_{\substack{\ell=1 \\ \ell \neq k}}^m \sum_{z \in \mathcal{P}_\ell} a_{jz}, & \text{if } \ell = k \\ \sum_{z \in \mathcal{P}_\ell} a_{jz}, & \text{otherwise.} \end{cases}$$

To conclude, $\|G^{(k)}\| \leq \sqrt{n_{\text{intra},k}} \max_{(i,j) \in \mathcal{E}_{\text{span},k}} |G_{ij}^{(k)}|$, and due to (A3), $\tilde{\gamma}_{ij}^{(k\ell)} = \tilde{\gamma}^{(k\ell)}$ is independent of i and j . Thus, $\|G^{(k)}\| \leq \sum_{\ell=1}^m 2 \max_r n_{\text{intra},r} \tilde{\gamma}^{(k\ell)} \|x_{\text{intra}}^{(\ell)}\|$, and this concludes the proof. ■

Proof of Theorem 3.2: The system (8) can be viewed as the perturbation via $G(x_{\text{intra}}, x_{\text{inter}})$ of $\dot{x}_{\text{intra}} = F(x_{\text{intra}})$, which describes the dynamics of m disjoint networks of oscillators

$$\dot{x}_{\text{intra}}^{(k)} = F^{(k)}(x_{\text{intra}}^{(k)}). \quad (21)$$

The origin of each system (21) is an exponentially stable equilibrium, which can be shown with the Lyapunov candidate

$$V_k(x_{\text{intra}}) = x_{\text{intra}}^{(k)\top} P_k x_{\text{intra}}^{(k)}$$

where $P_k \succ 0$ is such that $J_k^\top P_k + P_k J_k = -Q_k$ for $Q_k \succ 0$. In fact, the derivative of V along the trajectories (21) is

$$\begin{aligned} \dot{V}_k(x_{\text{intra}}^{(k)}) = & F^{(k)\top}(x_{\text{intra}}^{(k)}) P_k x_{\text{intra}}^{(k)} + x_{\text{intra}}^{(k)\top} P_k F^{(k)}(x_{\text{intra}}^{(k)}) \\ = & x_{\text{intra}}^{(k)\top} (J_k^\top P_k + P_k J_k) x_{\text{intra}}^{(k)} + O(\|x_{\text{intra}}^{(k)}\|^3) \quad (22) \end{aligned}$$

¹⁰Letting $\delta = \sin(a + b) - \sin(a)$, we have $|\delta| = |2 \sin(\frac{b}{2}) \cos(a + \frac{b}{2})| \leq |2 \sin(\frac{b}{2})|$, from which the inequality $|\delta| \leq |b|$ follows.

and the latter is strictly negative when $\|x_{\text{intra}}^{(k)}\| \leq r$ and $r \in \mathbb{R}_{>0}$ is sufficiently small. Further, it holds that: 1) $\|\partial V_k / \partial x_{\text{intra}}^{(k)}\| \leq 2\lambda_{\max}(P_k)\|x_{\text{intra}}^{(k)}\|$; 2) $\dot{V}_k(x_{\text{intra}}^{(k)}) \leq -\lambda_{\min}(Q_k)\|x_{\text{intra}}^{(k)}\|^2$; and 3) the perturbation terms $G^{(k)}(x_{\text{intra}}, x_{\text{inter}})$ are linearly bounded in $\|x_{\text{intra}}^{(k)}\|$ following statement 2) in Lemma 3.1.

Consider now the following Lyapunov candidate for (8):

$$V(x_{\text{intra}}) = \sum_{k=1}^m d_k V_k(x_{\text{intra}}^{(k)}), \quad d_k > 0.$$

From [54, ch. 9.5], we have

$$\dot{V}(x_{\text{intra}}) \leq -\frac{1}{2}(DS + S^T D)\|x_{\text{intra}}\|^2 \quad (23)$$

where $D = \text{diag}(d_1, \dots, d_m)$, and S satisfies

$$S = [s_{k\ell}] = \begin{cases} \frac{\lambda_{\min}(Q_k)}{\lambda_{\max}(P_k)} - \gamma^{(kk)}, & \text{if } k = \ell \\ -\gamma^{(k\ell)}, & \text{if } k \neq \ell. \end{cases} \quad (24)$$

The origin of (8) is locally exponentially stable if S is an M -matrix [54, Lemma 9.7 and Th. 9.2]. Finally, choosing $Q_k = I$ in (24) yields condition (14) in Theorem 3.2. ■

B. Proofs of the Results in Section III-B

Let \mathcal{C} be the set of connected clusters pairs, that is,

$$\mathcal{C} = \{(\ell, z) : \exists (i, j) \in \mathcal{E} \text{ with } i \in \mathcal{P}_\ell, j \in \mathcal{P}_z, \text{ and } \ell < z\}.$$

With a slight abuse of notation, for any $(\ell, z) \in \mathcal{C}$, we define $x^{(\ell z)} = x_{ij}$, for any node $i \in \mathcal{P}_\ell$ and $j \in \mathcal{P}_z$.

Lemma A.1. (Linearized intracluster dynamics): The linearization of the intracluster dynamics (8) around the trajectory $x_{\text{intra}} = 0$ and $x_{\text{inter}} = x_{\text{nom}}$ reads as follows:

$$\dot{x}_{\text{intra}} = (J_{\text{intra}} + J_{\text{inter}}) x_{\text{intra}} \quad (25)$$

where J_{intra} is defined in Lemma 3.1, and

$$J_{\text{inter}} = \frac{\partial G}{\partial x_{\text{intra}}} \Big|_{\substack{x_{\text{intra}}=0 \\ x_{\text{inter}}=x_{\text{nom}}}} \triangleq \sum_{(\ell, z) \in \mathcal{C}} \cos(x^{(\ell z)}) J_{\text{inter}}^{(\ell z)}.$$

Proof: Linearization of (8) around the trajectory $(x_{\text{intra}}, x_{\text{inter}}) = (0, x_{\text{nom}})$ yields $\partial F / \partial x_{\text{intra}} = J_{\text{intra}}$ and $\partial G / \partial x_{\text{intra}} = J_{\text{inter}}$. The remaining derivatives vanish. That is, $\partial F / \partial x_{\text{inter}} = 0$ because F does not depend on x_{inter} , and $\partial G / \partial x_{\text{inter}} = 0$ because of Assumption (A3). In fact, for any intracluster difference x_{ij} with $i, j \in \mathcal{P}_\ell$, $\ell \in \{1, \dots, m\}$

$$\frac{\partial G_{ij}}{\partial x_{\text{inter}}} \Big|_{\substack{x_{\text{intra}}=0 \\ x_{\text{inter}}=x_{\text{nom}}}} = \sum_{(\ell, z) \in \mathcal{C}} \cos(x^{(\ell z)}) \underbrace{\sum_{k \in \mathcal{P}_z} [a_{jk} - a_{ik}]}_{=0} = 0.$$

This concludes the proof. ■

We next characterize an asymptotic property of the intercluster differences through the following instrumental result.

Lemma A.2. (Asymptotic behavior of the intercluster dynamics for large frequency differences): Let $i \in \mathcal{P}_\ell$, $j \in \mathcal{P}_z$, and $\ell \neq z$. Then, the intercluster difference x_{ij} satisfies

$$\lim_{|\omega_j - \omega_i| \rightarrow \infty} \frac{x_{ij}(t)}{\omega_j - \omega_i} = t. \quad (26)$$

Proof: Let $\bar{\omega}_{ij} = \omega_j - \omega_i$. We rewrite (4) as

$$\begin{aligned} \dot{x}_{ij} &= \bar{\omega}_{ij} - (a_{ij} + a_{ji}) \sin(x_{ij}) \\ &\quad + \sum_{k \neq i, j} [a_{jk} \sin(x_{jk}) - a_{ik} \sin(x_{ik})]. \end{aligned} \quad (27)$$

From (27), let $\beta = \sum_{k \neq i, j} [a_{jk} + a_{ik}]$, and

$$\dot{\underline{x}}_{ij} = \bar{\omega}_{ij} - (a_{ij} + a_{ji}) \sin(\underline{x}_{ij}) - \beta \quad (28)$$

$$\dot{\bar{x}}_{ij} = \bar{\omega}_{ij} - (a_{ij} + a_{ji}) \sin(\bar{x}_{ij}) + \beta \quad (29)$$

with $\underline{x}_{ij}(0) = \bar{x}_{ij}(0) = x_{ij}(0)$. Integrating (28) yields

$$\int_{x_{ij}(0)}^{\underline{x}_{ij}(t)} \frac{dy}{\bar{\omega}_{ij} - (a_{ij} + a_{ji}) \sin(y) - \beta} = \int_0^t d\tau. \quad (30)$$

As $|\bar{\omega}_{ij}|$ grows, it holds that $|(a_{ij} + a_{ji}) + \beta| < |\bar{\omega}_{ij}|$. Therefore

$$\begin{aligned} \frac{1}{\bar{\omega}_{ij} - (a_{ij} + a_{ji}) \sin(y) - \beta} &= \frac{1}{\bar{\omega}_{ij}} \left[\frac{1}{1 - \frac{(a_{ij} + a_{ji}) \sin(y) + \beta}{\bar{\omega}_{ij}}} \right] \\ &= \frac{1}{\bar{\omega}_{ij}} \sum_{k=0}^{\infty} \left[\frac{(a_{ij} + a_{ji}) \sin(y) + \beta}{\bar{\omega}_{ij}} \right]^k. \end{aligned}$$

In view of the latter equality, (30) becomes

$$\begin{aligned} t &= \frac{\underline{x}_{ij}(t) - x_{ij}(0)}{\bar{\omega}_{ij}} \\ &\quad + \frac{1}{\bar{\omega}_{ij}} \underbrace{\int_{x_{ij}(0)}^{\underline{x}_{ij}(t)} \sum_{k=1}^{\infty} \left[\frac{(a_{ij} + a_{ji}) \sin(y) + \beta}{\bar{\omega}_{ij}} \right]^k dy}_{O(\bar{\omega}_{ij}^{-1})} \end{aligned}$$

or, equivalently

$$\underline{x}_{ij}(t) = \bar{\omega}_{ij} t + x_{ij}(0) + O(\bar{\omega}_{ij}^{-1}). \quad (31)$$

Similarly, the solution of (29) has the form in (31). Finally, using the comparison principle [54, Lemma 3.4], it holds that $\underline{x}_{ij}(t) \leq x_{ij}(t) \leq \bar{x}_{ij}(t)$ for all $t \geq 0$. Hence, $\frac{x_{ij}(t)}{\bar{\omega}_{ij}} \rightarrow t$ as $|\bar{\omega}_{ij}| \rightarrow \infty$ and this concludes the proof. ■

We are now ready to prove Theorem 3.3.

Proof of Theorem 3.3: Consider the Lyapunov candidate $V(x_{\text{intra}}, t) = x_{\text{intra}}^T \Gamma(t) x_{\text{intra}}$, and notice that, using (25),

$$\begin{aligned} \dot{V}(x_{\text{intra}}, t) &= \dot{x}_{\text{intra}}^T \Gamma x_{\text{intra}} + x_{\text{intra}}^T \dot{\Gamma} x_{\text{intra}} + x_{\text{intra}}^T \Gamma \dot{x}_{\text{intra}} \\ &= x_{\text{intra}}^T \left[J_{\text{intra}}^T \Gamma + \Gamma J_{\text{intra}} + \dot{\Gamma} \right. \\ &\quad \left. + \sum_{(\ell, z) \in \mathcal{C}} \cos(x^{(\ell z)}) \left(J_{\text{inter}}^{(\ell z)T} \Gamma + \Gamma J_{\text{inter}}^{(\ell z)} \right) \right] x_{\text{intra}} + O(\|x_{\text{intra}}\|^3). \end{aligned} \quad (32)$$

Let

$$\dot{\Gamma} = - \sum_{(\ell, z) \in \mathcal{C}} \cos(x^{(\ell z)}) \left(J_{\text{inter}}^{(\ell z)T} \Gamma + \Gamma J_{\text{inter}}^{(\ell z)} \right). \quad (33)$$

When the intercluster natural frequencies satisfy $|\omega_i - \omega_j| \rightarrow \infty$ for all i, j , then $\Gamma(t) \rightarrow \Gamma(0)$ for all times t . In fact, integrating both sides of (33) and substituting $\Gamma(t) = \Gamma(0)$ yields

$$\begin{aligned} \int_0^t \dot{\Gamma} d\tau &= \Gamma(t) - \Gamma(0) = \Gamma(0) - \Gamma(0) = 0 \\ &= - \sum_{(\ell, z) \in \mathcal{C}} \int_0^t \cos(x^{(\ell z)}) \left(J_{\text{inter}}^{(\ell z)T} \Gamma + \Gamma J_{\text{inter}}^{(\ell z)} \right) d\tau \\ &= - \sum_{(\ell, z) \in \mathcal{C}} \left(J_{\text{inter}}^{(\ell z)T} \Gamma(0) + \Gamma(0) J_{\text{inter}}^{(\ell z)} \right) \int_0^t \cos(x^{(\ell z)}) d\tau \end{aligned}$$

which holds true because $\int \cos(x^{(\ell z)}) d\tau = 0$ due to Lemma A.2. Because J_{intra} is stable, we conclude that, when the intercluster natural frequencies satisfy $|\omega_i - \omega_j| \rightarrow \infty$ for all i, j , $\dot{\Gamma} = 0$, and there exists $\Gamma(0)$ such that (32) is strictly negative. This concludes the proof of the claimed statement. ■

Proof of Lemma 3.4: When $x_{\text{intra}} = 0$, the differential equation (27) reduces to $\dot{x}_{\text{inter}} = \bar{\omega} - \bar{a} \sin(x_{\text{inter}})$, which is a separable differential equation with solution as in (15). To show that the period of (15) is equal to $T = 2\pi/\sqrt{\bar{\omega}^2 - \bar{a}^2}$, we assume, without loss of generality, that $\tau = 0$. It is easy to see that, because $\tan(t)$ is π -periodic, $x_{\text{nom}}(t) = x_{\text{nom}}(t + 2\pi/\sqrt{\bar{\omega}^2 - \bar{a}^2})$. Further, notice that the variable substitution $z = x_{\text{nom}}$ in $\int_0^t \cos(x_{\text{nom}}) d\tau$ yields

$$\begin{aligned} \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau &= \int_{x_{\text{nom}}(0)}^{x_{\text{nom}}(t)} \frac{\cos(z)}{\bar{\omega} - \bar{a} \sin(z)} dz \\ &= \frac{1}{\bar{a}} \log \left(\frac{\bar{\omega} - \bar{a} \sin(x(0))}{\bar{\omega} - \bar{a} \sin(x_{\text{nom}}(t))} \right) \end{aligned} \quad (34)$$

which implies the bound (16). To prove that $\cos(x_{\text{nom}})$ has zero time average, it suffices to substitute $t = T$ into (34). ■

Proof of Theorem 3.5: Consider the Lyapunov candidate $V(x_{\text{intra}}, t) = x_{\text{intra}}^T \Gamma(t) x_{\text{intra}}$, and notice that, using (25)

$$\begin{aligned} \dot{V}(x_{\text{intra}}, t) &= x_{\text{intra}}^T [J_{\text{intra}}^T \Gamma + \Gamma J_{\text{intra}} + \dot{\Gamma} + \cos(x_{\text{nom}}) \\ &\quad (J_{\text{inter}}^T \Gamma + \Gamma J_{\text{inter}})] x_{\text{intra}} + O(\|x_{\text{intra}}\|^3). \end{aligned} \quad (35)$$

Let $\dot{\Gamma} = -\cos(x_{\text{nom}})(J_{\text{inter}}^T \Gamma + \Gamma J_{\text{inter}})$ and notice that, following [57, Exercise 3.9 and Property 4.2], its solution satisfies

$$\begin{aligned} \Gamma(t) &= \exp \left[- \int_0^t \cos(x_{\text{nom}}(\tau)) J_{\text{inter}}^T d\tau \right] \Gamma(0) \\ &\quad \cdot \exp \left[- \int_0^t \cos(x_{\text{nom}}(\tau)) J_{\text{inter}} d\tau \right]. \end{aligned}$$

This implies that $V(x_{\text{intra}}, t)$ is a Lyapunov function for (25) because, by Lemma 3.4, $\int_0^t \cos(x_{\text{nom}}(\tau)) d\tau$ is bounded. Furthermore, notice that

$$\begin{aligned} &\exp \left[- \int_0^t \cos(x_{\text{nom}}(\tau)) J_{\text{inter}}^T d\tau \right] \\ &= I + \underbrace{\sum_{k=1}^{\infty} \frac{(J_{\text{inter}}^T)^k}{k!} \left(- \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau \right)^k}_{\Delta}. \end{aligned}$$

Thus, (35) can equivalently be written as $\dot{V} = x_{\text{intra}}^T [J_{\text{intra}}^T \Gamma(0) + \Gamma(0) J_{\text{intra}} + M] x_{\text{intra}} + O(\|x_{\text{intra}}\|^3)$, where $M = J_{\text{intra}}^T \Delta \Gamma(0) \Delta^T + \Delta \Gamma(0) \Delta^T J_{\text{intra}} + J_{\text{intra}}^T (\Delta \Gamma(0) + \Gamma(0) \Delta) + (\Delta \Gamma(0) + \Gamma(0) \Delta) J_{\text{intra}}$. Using the triangle inequality and Lemma 3.4, we obtain

$$\begin{aligned} \|\Delta\| &= \left\| \sum_{k=1}^{\infty} \frac{(J_{\text{inter}}^T)^k}{k!} \left(- \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau \right)^k \right\| \\ &\leq \sum_{k=1}^{\infty} \frac{\|J_{\text{inter}}\|^k}{k!} \left| \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau \right|^k \\ &= e^{\left| \int_0^t \cos(x_{\text{nom}}(\tau)) d\tau \right| \|J_{\text{inter}}\|} - 1 \leq e^{\frac{1}{\bar{a}} \log \left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right) \|J_{\text{inter}}\|} - 1. \end{aligned}$$

Because J_{intra} is stable, there always exists $\Gamma(0) \succ 0$ such that $J_{\text{intra}}^T \Gamma(0) + \Gamma(0) J_{\text{intra}} = -Q$ for any $Q \succ 0$. Thus

$$\dot{V} \leq (-\lambda_{\min}(Q) + \|M\|) \|x_{\text{intra}}\|^2 + O(\|x_{\text{intra}}\|^3). \quad (36)$$

By a simple Lyapunov argument, the cluster synchronization manifold \mathcal{S}_p is locally exponentially stable if $\|M\| < \lambda_{\min}(Q)$. In addition, $\|M\|$ can be upper bounded as $\|M\| \leq 2\|J_{\text{intra}}\| \|\Gamma(0)\| \|\Delta\| (\|\Delta\| + 2) \leq 2\lambda_{\max}(\Gamma(0)) \|J_{\text{intra}}\| (e^{\frac{2}{\bar{a}} \log \left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right) \|J_{\text{inter}}\|} - 1)$. Thus, a sufficient condition for local exponential stability is $2\lambda_{\max}(\Gamma(0)) \|J_{\text{intra}}\| (e^{\frac{2}{\bar{a}} \log \left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right) \|J_{\text{inter}}\|} - 1) < \lambda_{\min}(Q)$, and because the ratio $\lambda_{\min}(Q)/\lambda_{\max}(\Gamma(0))$ is maximized for $Q = I$ [54, Exercise 9.1], we have

$$2\lambda_{\max}(\Gamma(0)) \|J_{\text{intra}}\| \left(e^{\frac{2}{\bar{a}} \log \left(\frac{\bar{\omega} + \bar{a}}{\bar{\omega} - \bar{a}} \right) \|J_{\text{inter}}\|} - 1 \right) < 1$$

from which condition (18) follows. ■

Proof of Theorem 3.6: From (35) and for $\beta \in \mathbb{R}_{>0}$, we have $\dot{V}(x_{\text{intra}}, t) = x_{\text{intra}}^T [J_{\text{intra}}^T \Gamma + \Gamma J_{\text{intra}}] x_{\text{intra}} + O(\|x_{\text{intra}}\|^3) = -\beta x_{\text{intra}}^T \Gamma x_{\text{intra}} + O(\|x_{\text{intra}}\|^3)$, which is negative in a small neighborhood of the origin. ■

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Tommaso Menara (S'17) received the Laurea degree (B.Sc. equivalent) in mechatronics engineering from the University of Padova, Padova, Italy, in 2013, and the Laurea Magistrale degree (M.Sc. equivalent) in robotics and automation engineering from the University of Pisa, Pisa, Italy, in 2016. He is currently working toward the Ph.D. degree with the Department of Mechanical Engineering, University of California at Riverside, Riverside, CA, USA.

His research interests include control of complex networks and network neuroscience.



Giacomo Baggio received the Ph.D. degree in control systems engineering from the University of Padova, Padova, Italy, in 2018.

He is currently a Postdoctoral Scholar with the Department of Mechanical Engineering, University of California at Riverside, Riverside, CA, USA. From 2015 to 2016, he was a Visiting Scholar with the Department of Engineering, University of Cambridge, Cambridge, U.K. His current research interests are in the analysis and control of dynamical networks.

Dr. Baggio was the recipient of the Best Student Paper Award at the 2018 European Control Conference.



Danielle S. Bassett received the B.S. degree from Penn State University, University Park, PA, USA, in 2004, and the CPGS and Ph.D. degrees from the University of Cambridge, Cambridge, U.K., as a Churchill Scholar and as an NIH Health Sciences Scholar, in 2005 and 2009, respectively, all in physics.

She is the Eduardo D. Glandt Faculty Fellow and an Associate Professor with the Department of Bioengineering, University of Pennsylvania, Philadelphia, PA, USA. She is most well known

for her work blending neural and systems engineering to identify fundamental mechanisms of cognition and disease in human brain networks. She was a Postdoctoral Researcher with the UC Santa Barbara and a Junior Research Fellow with the Sage Center for the Study of the Mind. She is the author of more than 200 peer-reviewed publications, which have garnered more than 15900 citations, as well as numerous book chapters and teaching materials.

Dr. Bassett was the recipient of multiple prestigious awards, including American Psychological Association's "Rising Star" (2012), Alfred P Sloan Research Fellow (2014), MacArthur Fellow Genius Grant (2014), Early Academic Achievement Award from the IEEE Engineering in Medicine and Biology Society (2015), Harvard Higher Education Leader (2015), Office of Naval Research Young Investigator (2015), National Science Foundation CAREER (2016), Popular Science Brilliant 10 (2016), Lagrange Prize in Complex Systems Science (2017), and Erdős-Rényi Prize in Network Science (2018).



Fabio Pasqualetti (S'07–M'13) received the Laurea degree (B.Sc. equivalent) in computer engineering and the Laurea Magistrale degree (M.Sc. equivalent) in automation engineering from the University of Pisa, Pisa, Italy, in 2004 and 2007, respectively, and the Doctor of Philosophy degree in mechanical engineering from the University of California, Santa Barbara, CA, USA, in 2012.

He is an Assistant Professor with the Department of Mechanical Engineering, University of California at Riverside, Riverside, CA. His main research interests include the analysis and control of complex networks, security of cyber-physical systems, distributed control, and network neuroscience.

Dr. Pasqualetti was the recipient of several awards, including a Young Investigator Program award from the Army Research Office in 2017 and the 2016 Transactions on Control of Network Systems Outstanding Paper Award from the IEEE Control Systems Society.